

Superpotentials for quiver gauge theories

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ABSTRACT: We compute superpotentials for quiver gauge theories arising from marginal D-Brane decay on collapsed del Pezzo cycles S in a Calabi-Yau X . This is done using the machinery of A_∞ products in the derived category of coherent sheaves of X , which in turn is related to the derived category of S and quiver path algebras. We confirm that the superpotential is what one might have guessed from analyzing the moduli space, i.e., it is linear in the fields corresponding to the Ext^2 s of the quiver and that each such Ext^2 multiplies a polynomial in Ext^1 s equal to precisely the relation represented by the Ext^2 .

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1. Introduction

Singularities of string backgrounds have attracted much attention and have been investigated using a variety of methods [1–6]. One is to study the gauge theory on a D -brane probe of the singularity. While there has been much work done on extracting gauge theory data for various types of singularities (abelian [1] and non-abelian [7] orbifolds, conifolds [8, 9], toric [10–12, 6] and generalized del Pezzo [13, 14] singularities), a general method for determining the superpotential has been lacking. In [10] the superpotential was obtained from integrating the quiver relations for certain cases, with various ad-hoc methods being used to resolve ambiguities that arise in such an integration. In this paper, using previous work of [15], we present a general rigorous method for obtaining the superpotential of such quiver gauge theories from the quiver relations. We show that the superpotential is just the naive sum of terms of the form relation times the Ext^2 field corresponding to the relation. We apply the method to the trivial example of a \mathbb{P}^2 as well as to a $d\mathbb{P}_1$, in which case we

get a non-homogeneous superpotential. In principle we can apply it to a general del Pezzo singularity — all we need is the del Pezzo quiver and relations.

We deal with flat compactifications $M \times X$ where M is 4 dimensional Minkowski space and X is a Calabi-Yau manifold, and we probe the theory with space filling branes — i.e., $(n+3)$ -branes, where n is the dimensionality of the brane within X . We expect such branes to be BPS and stable when probing a smooth point of X , but to marginally decay into a collection of so-called fractional branes when the point becomes singular. We consider singularities obtained when a complex surface (i.e., one that has 4 real dimensions) S shrinks down to zero size in X by varying the Kähler parameters. Assuming that S is smooth and irreducible, it is known that S must be a del Pezzo surface, i.e., $\mathbb{P}^1 \times \mathbb{P}^1$, or \mathbb{P}^2 blown up at m points (denoted by dP_m), where m ranges from 0 to 8. The $3+1$ dimensional quiver gauge theory associated to this marginal decay into fractional branes is the one whose superpotential we are after. This set up has been studied extensively in the literature [16, 9, 17–21, 14, 22, 23].

The moduli space of a D-brane is given by the space of critical points of the superpotential. Thus, knowing the moduli space one may make a guess at the form of the superpotential. In the case at hand, this leads to a natural conjecture for the superpotential. By using more rigorous methods we are able to show that this conjecture is correct.

For our analysis we will use the algebraic machinery of the derived category of coherent sheaves, developed in particular in [24, 25, 5]. The fractional branes have tractable representations as elements of the derived category $\mathbf{D}(X)$ of coherent sheaves on X . In principle the method of [15] can be applied to find the so-called A_∞ products in the algebra of Ext groups. These A_∞ products are determined by combinatorial relations that they have to satisfy (coming from Feynman diagrams in the associated topological theory) and they encode the superpotential. Applying the technique of [15] directly is difficult and in order to make the problem tractable, we instead proceed in two steps. First, we use a spectral sequence argument to reduce the problem from one of studying sheaves on X to the simpler one of studying sheaves on S . It is in this reduction that we show that each Ext^2 appears linearly in the superpotential, multiplying a term that involves only Ext^1 s and is determined by A_∞ relations over S . To compute these we exploit the well understood properties of $\mathbf{D}(S)$, and specifically its intimate relation with the derived category of quiver representations [26, 27]. We see that the terms involving the Ext^1 s are just the (possibly non-homogeneous) relations in the quiver. It is important to note that we obtain the superpotential only up to certain nonlinear field redefinitions (see [15]) — this is the most that could be expected from such topological sigma model methods as we use.

The plan of the paper is as follows: in section 2 we review quivers and sheaves on del Pezzo surfaces, and how they relate. In section 3 we review A_∞ algebras and the method of [15] for using the topological B model to compute A_∞ structure and hence the superpotential. In section 4 we introduce the quiver gauge theory we want to study and prove that its superpotential is linear in the Ext^2 s, which multiply terms determined by the A_∞ structure over S — this is the reduction from sheaves on X to sheaves on S . In section 5 we solve the problem on S by reframing it as a computation in the derived category of quiver representations, and apply the solution to the case $S = \mathbb{P}^2$ and the more nontrivial

case $S = dP_1$. This example illustrates the general algorithm that can be carried through for any quiver with known relations.

2. Quivers and Sheaves

2.1 Quivers

We now review the necessary mathematical notions relating to quivers and their representations. Further background can be found in [5]. First, a quiver is a directed graph Q that consists of nodes v_i and arrows a_α . Its path algebra A is defined as follows: as a vector space, A is generated by all of the paths constructed through concatenation of arrows in Q . The product structure of A is defined on these generators as follows: if path 1 ends on the same node that path 2 begins on, the product is defined to be the obvious concatenation; otherwise it is defined to be 0. Note that corresponding to each node v_i we have a corresponding zero length path and hence an idempotent element e_i of A .

In the remainder of the paper, we will deal with a slight generalization to a quiver with relations. This is just a quiver whose path algebra is defined as the above A quotiented out by a subspace generated by linear combinations of paths called relations. We stipulate that any given relation must be a linear combination of paths between the same two nodes. It does not, however, need to be homogeneous. A simple example (with homogeneous relations) is the so-called Beilinson quiver, defined as:

$$\begin{array}{ccccc}
 & & a_0 & & b_0 \\
 & \circ & \leftarrow & \circ & \leftarrow & \circ \\
 & \swarrow & a_1 & \swarrow & b_1 & \swarrow \\
 & & & & & \\
 & \searrow & a_2 & \searrow & b_2 & \searrow \\
 & \circ & & \circ & & \circ \\
 v_0 & & & v_1 & & v_2
 \end{array} \tag{2.1}$$

with relations $a_\alpha b_\beta - a_\beta b_\alpha$.

For a given quiver Q , we can consider the associated category $A\text{-mod}$ of left A -modules. For any left A -module V we can form the vector spaces $V_i = e_i V$, of dimension N_i ; if we then think of each V_i as living on node v_i then we see that multiplication by any arrow a_α acts as a linear transformation between the spaces V_i at the tail and head of a_α , and these linear transformations respect the relations. This structure is known as a representation of the quiver Q , of dimension (N_i) .

A map ϕ between left A -modules V and W is simply a linear transformation that commutes with the action of A , i.e., $\phi(av) = a\phi(v)$. If we think of V and W as quiver representations then this condition is just the obvious constraint that the maps from V_i to W_i must commute with the linear transformations induced by the arrows, i.e., ϕ is a map of representations. Thus we sometimes refer to the category of left A -modules as the category of quiver representations, and we use these terms interchangeably from now on.

Corresponding to each node of Q there are two distinguished representations P_i and L_i ; when we make the connection to sheaves on del Pezzo surfaces below, these will correspond to sheaves in the exceptional collection and fractional branes, respectively, as we will see. L_i is defined simply as the one dimensional representation with $V_i = \mathbb{C}$ and all other $V_j = 0$. P_i is defined as the subspace of A generated by all paths that begin at v_i ; it is trivially seen to be a subrepresentation. It may seem that P_i is a rather large representation, and indeed, if the quiver has any loops there will be infinite dimensional P_i . However, in the case of

quivers associated to del Pezzo's there will be an ordering on the nodes that is respected by the arrows and hence all P_i will be finite-dimensional. This finite dimensionality is in large part responsible for the tractability of the problem and motivates the reduction mentioned earlier from sheaves on the Calabi-Yau to sheaves on the del Pezzo. In the Beilinson quiver, for example, we see that P_0 , P_1 , and P_2 have dimensions $(1, 0, 0)$, $(3, 1, 0)$, and $(6, 3, 1)$ respectively.

When we make the connection between quivers and sheaves it will be through the derived category. Before we talk about that, however, let us first discuss some basic homological properties of quivers. First of all, one can show that the P_i are projective objects in $A\text{-mod}$. In fact, they form a complete set in the sense that any left A -module has a resolution by various direct sums of these P_i . These projective resolutions can be used to compute higher Ext groups. For example, the projective resolution of any L_i is

$$\dots \longrightarrow \bigoplus_k P_k^{\oplus r_{ik}} \longrightarrow \bigoplus_k P_k^{\oplus n_{ik}} \longrightarrow P_i \longrightarrow L_i \longrightarrow 0. \tag{2.2}$$

Here n_{ij} is the number of arrows from node i to node j and r_{ij} is the number of independent relations imposed on paths from i to j . In the case of the Beilinson quiver, the resolutions are:

$$\begin{aligned} 0 &\longrightarrow P_0 \longrightarrow L_0 \longrightarrow 0 \\ 0 &\longrightarrow P_0^{\oplus 3} \longrightarrow P_1 \longrightarrow L_1 \longrightarrow 0 \\ 0 &\longrightarrow P_0^{\oplus 3} \longrightarrow P_1^{\oplus 3} \longrightarrow P_2 \longrightarrow L_2 \longrightarrow 0 \end{aligned} \tag{2.3}$$

Noting that Ext^k is the k th derived functor of Hom and that $\text{Hom}(P_i, L_j) = \delta_{ij}\mathbb{C}$ we can compute $\text{Ext}^p(A, B)$ by taking a projective resolution

$$\dots \longrightarrow \Pi_2 \longrightarrow \Pi_1 \longrightarrow \Pi_0 \longrightarrow A \longrightarrow 0, \tag{2.4}$$

where the Π_i are direct sums of P_i 's and from it constructing the complex

$$0 \longrightarrow \text{Hom}(\Pi_0, B) \longrightarrow \text{Hom}(\Pi_1, B) \longrightarrow \text{Hom}(\Pi_2, B) \longrightarrow \dots \tag{2.5}$$

The cohomology of this complex in the p th position is then $\text{Ext}^p(A, B)$. Using this method one can show that

$$\begin{aligned} \dim \text{Ext}^1(L_i, L_j) &= n_{ij} \\ \dim \text{Ext}^2(L_i, L_j) &= r_{ij}. \end{aligned} \tag{2.6}$$

In general, higher Ext's may also exist. For example, Ext^3 represents relations amongst relations. However, for the purposes of quiver gauge theories, the appearance of higher Ext's is unphysical [16, 28] and so we assume

$$\text{Ext}^k(L_i, L_j) = 0, \quad k \geq 3. \tag{2.7}$$

Finally, we note that we can recover the quiver path algebra from the projective representations P_i . Supposing Q is a quiver with n nodes, this is done as follows: we let

$$T = P_0 \oplus P_1 \oplus \cdots \oplus P_{n-1}. \tag{2.8}$$

Using the fact that $\text{Hom}(P_i, P_j)$ is simply the vector space of paths from j to i we can verify that

$$A \cong \text{End}(T)^{\text{op}}. \tag{2.9}$$

In other words, A is just the algebra $\text{End}(T)$ with the product structure reversed.

2.2 The Derived Category

Having reviewed this preliminary material about quiver representations we move on to briefly discuss the derived category. As mentioned above, the derived category will form a bridge between the quiver representations that we have already discussed and the category of coherent sheaves introduced below. One source is [5]; here we just review the facts we will use.

Given any abelian category \mathcal{A} (such as that of quiver representations, or that of coherent sheaves discussed below) we can define its derived category $\mathbf{D}(\mathcal{A})$ as follows. The objects in $\mathbf{D}(\mathcal{A})$ are complexes of objects in \mathcal{A} :

$$\cdots \longrightarrow \mathcal{E}^0 \xrightarrow{d_0} \mathcal{E}^1 \xrightarrow{d_1} \mathcal{E}^2 \xrightarrow{d_2} \cdots, \tag{2.10}$$

To construct the morphisms in $\mathbf{D}(\mathcal{A})$, we begin with the abelian group of all possible maps between complexes (not necessarily respecting the differential). These maps are graded by their degree p and can be written as

$$\sum_n f_{n,n+p} \tag{2.11}$$

where $f_{m,n}$ is a map from \mathcal{E}^m to \mathcal{E}^n . We define a differential on this group by (abusing notation slightly):

$$(df)_{n,p+1} = d_{n+p}f_{n,p} - (-1)^p f_{n+1,p}d_n \tag{2.12}$$

The derived morphisms are now defined as the cohomology of this group, with formal inverses added in for all quasi-isomorphisms (that is, those chain maps which induce isomorphisms on cohomology).

Now we state some necessary results without proof. Given any object A in \mathcal{A} , we can construct the associated one term complex whose only nonzero entry is A , at the zeroth position. For brevity we will henceforth refer to both the object and the associated one term complex by A . Then, for A, B in \mathcal{A} , $\text{Ext}^\bullet(A, B)$ is given by the group of derived morphisms between the complexes associated to A and B , with the grading on Ext corresponding to the grading of the derived morphisms. In fact, this is the generalization of the notion of Ext to the arbitrary elements of $\mathbf{D}(\mathcal{A})$. Also, any A in \mathcal{A} is equivalent to its projective

or injective resolution in the derived category. Further, if we represent either A by its projective resolution or B by its injective resolution then the generators of $\text{Ext}(A, B)$ can be written as honest chain maps between these complexes.

2.3 Sheaves

We now turn to reviewing key aspects of the other relevant category, that of coherent sheaves. It turns out that (as we will see in more detail below) the derived category of coherent sheaves on a Calabi-Yau manifold X , denoted $\mathbf{D}(X)$, precisely describes D -branes in the topological B model defined on X . The open string modes stretching between them are described by the Ext groups of the sheaf homs between the relevant branes [24, 25, 5]. These, in turn, describe the massless spectrum of the physical theory on $M \times X$. In fact $\mathbf{D}(X)$ contains enough information to determine the tree-level superpotential of the low energy effective theory, in the form of A_∞ products. We will discuss all of this below, but for now let us start by introducing coherent sheaves on Calabi-Yau's and del Pezzo's.

The category of coherent sheaves on a space X is an enlargement of the category of vector bundles (also referred to as “locally free sheaves”) on X — it contains vector bundles as well as all kernels and cokernels of maps of vector bundles. For a precise definition, starting from the general concept of a sheaf, see [5] or [29]. We can very roughly think of it as including, in addition to vector bundles over X , more exotic objects such as vector bundles over submanifolds of X .

In the physical problem we consider D -branes on a shrinking cycle S which is embedded in X : $i : S \rightarrow X$. i induces an embedding $i_* : \mathbf{D}(S) \rightarrow \mathbf{D}(X)$, and it is no surprise that the branes we'll be interested in are in fact in the image of i_* . Now, $\mathbf{D}(S)$ has been studied extensively by mathematicians and is well understood.

We proceed by first defining a complete strongly exceptional collection of sheaves on S to be an ordered set $\{\mathcal{F}_0, \dots, \mathcal{F}_{n-1}\}$ that generates $\mathbf{D}(S)$ and satisfies $\text{Ext}_S^p(\mathcal{F}_i, \mathcal{F}_j) = 0$ for $p \neq 0$ and any i and j , and $\text{Ext}_S^0(\mathcal{F}_i, \mathcal{F}_j) = \text{Hom}_S(\mathcal{F}_i, \mathcal{F}_j) = 0$ for $i > j$ and $\text{Hom}_S(\mathcal{F}_i, \mathcal{F}_i) = \mathbb{C}$. Given such a complete strongly exceptional collection, we can define

$$A = \text{End}(\mathcal{F}_0 \oplus \mathcal{F}_1 \oplus \dots \oplus \mathcal{F}_{n-1})^{\text{op}} \tag{2.13}$$

It turns out that A is the path algebra of a quiver Q , and the \mathcal{F}_i are isomorphic (as A -modules) to the projective representations P_i defined above. Given this we can reconstruct the quiver uniquely simply by noting that $\text{Hom}_S(\mathcal{F}_i, \mathcal{F}_j) = \text{Hom}(P_i, P_j)$ is just the space of paths from node j to node i . In fact, Bondal [26] proved that the derived category of A -modules, $\mathbf{D}(A\text{-mod})$ is equivalent to $\mathbf{D}(S)$.

As an example, consider $S = \mathbb{P}^2$. An exceptional collection is given by $\{\mathcal{O}, \mathcal{O}(1), \mathcal{O}(2)\}$. We have $\text{Hom}(\mathcal{O}, \mathcal{O}(1)) \cong \mathbb{C}^3$, $\text{Hom}(\mathcal{O}(1), \mathcal{O}(2)) \cong \mathbb{C}^3$. Denote these maps, which are just multiplication by the homogeneous coordinates, by x_i and y_i respectively, $i = 1, 2, 3$. We also have $\text{Hom}(\mathcal{O}, \mathcal{O}(2)) \cong \mathbb{C}^6$ — these maps are multiplication by homogeneous degree two polynomials in the homogeneous coordinates. All this implies that we have three arrows x_i from node 2 to node 1, three arrows y_i from node 3 to node 2, and that all paths from node 3 to node 1 are compositions of these arrows, with relations $x_i y_j - x_j y_i = 0$.

Another example which will be thoroughly dealt with below is $S = \text{dP}_1$, which is \mathbb{P}^2 with one point blown up. Letting C_1 be the exceptional divisor, a complete strongly exceptional collection is $\{\mathcal{O}, \mathcal{O}(C_1), \mathcal{O}(H), \mathcal{O}(2H)\}$ where H is the hyperplane divisor. A slightly more involved analysis shows the quiver to be

$$(2.14)$$

with the relations $b_0d_1 - b_1d_0 = 0$, $ab_0d_2 - cd_0 = 0$, and $ab_1d_2 - cd_1 = 0$.

Now that we have Bondal’s theorem, we can use either $\mathbf{D}(S)$ or $\mathbf{D}(A\text{-mod})$ to describe branes on S . We will call the branes that correspond to the representations L_i fractional branes. Of course, we are actually interested in branes in X , i.e., in the image of $\mathbf{D}(S)$ in $\mathbf{D}(X)$ as noted above. Using a local model of the Calabi-Yau X , namely representing it as the total space of the normal bundle of S in X (which is isomorphic to the canonical bundle) one can determine the Ext groups of sheaves in $i_*\mathbf{D}(S)$ in terms of the Ext groups in $\mathbf{D}(S)$. Namely, we find using a spectral sequence argument [30, 31, 21, 32] that

$$\text{Ext}_X^p(i_*L_i, i_*L_j) = \text{Ext}_S^p(L_i, L_j) \oplus \text{Ext}_S^{3-p}(L_j, L_i). \quad (2.15)$$

In fact it is also true that only one of the direct summands on the right hand side of the above equation is nonzero, and so we see that embedding S in X creates new open string degrees of freedom — new Ext^1 ’s corresponding to reversing Ext^2 ’s in the del Pezzo quiver. We can add in arrows corresponding to these new Ext^1 ’s to obtain the completed quiver. For example, the completion of the Beilinson quiver is

$$(2.16)$$

while the completion of the dP_1 quiver becomes

$$(2.17)$$

3. Superpotentials from Topological Field Theory

3.1 Topological Field Theory

Having developed and reviewed the requisite mathematical machinery, let us get to the problem at hand, namely computing superpotentials for effective dimensionally reduced theories [15]. Our setting is, as we said, $M \times X$ with M being four dimensional Minkowski space and X a Calabi-Yau threefold. In general, the object is to figure out how to obtain the superpotential for a specified distribution of space-filling branes — the case of interest involves putting $D3$ branes (which look like points in X) on a collapsing del Pezzo cycle S in X , but let us for the purpose of developing some formalism first tackle the case of a single space-filling and Calabi-Yau filling $D9$ brane, described by a complex line bundle

$E \rightarrow X$ with a hermitian connection. (By itself this case is unphysical, in a sense, because of anomalies but the topological field theory makes perfect sense.)

In this case, the massless four dimensional field content is determined by the Dolbeault cohomology of X valued in $\text{End}(E)$, $H_{\bar{\partial}}^{0,q}(X, \text{End}(E))$. Specifically, the number of vector bosons is given by $H_{\bar{\partial}}^{0,0}(X, \text{End}(E)) = \text{End}(E)$, where by abuse of notation the second term refers to the space of global sections of $\text{End}(E)$. We will work with simple line bundles, for which $\text{End}(E) = \mathbb{C}$. We could of course also take N copies of the brane, $E^{\oplus N}$, whereby we obtain a $U(N)$ gauge boson. Likewise, the number of chiral superfields is given by $H_{\bar{\partial}}^{0,1}(X, \text{End}(E))$. Again, these are in the adjoint of $U(N)$ when we take the bundle to be $E^{\oplus N}$.

In order to get a term in the tree-level superpotential, we have to compute a disk diagram with boundary insertions of vertex operators that correspond to the chiral superfields that appear in that term. What makes this problem computationally tractable is the fact that this disk diagram can be computed in a topological theory [33]; it is in some sense protected from α' corrections. Specifically, the open string topological B -model on X with a D -brane E has open string spectrum given by $A = H_{\bar{\partial}}^{0,q}(X, \text{End}(E))$. Thus, if we define the disk correlation functions as:

$$B_{i_0, i_1, \dots, i_k} = (-1)^{\zeta_1 + \zeta_2 + \dots + \zeta_{k-1}} \langle \psi_{i_0} \psi_{i_1} P \int \psi_{i_2}^{(1)} \int \psi_{i_3}^{(1)} \dots \int \psi_{i_{k-1}}^{(1)} \psi_{i_k} \rangle, \quad (3.1)$$

Here the ψ_{i_m} are vertex operators of ghost number one, i.e., they correspond to states in $H_{\bar{\partial}}^{0,1}(X, \text{End}(E))$. If we let Z_i be the effective four dimensional superfield corresponding to the open string mode ψ_i , then the superpotential is

$$W = \text{Tr} \left(\sum_{k=2}^{\infty} \sum_{i_0, i_1, \dots, i_k} \frac{B_{i_0, i_1, \dots, i_k}}{k+1} Z_{i_0} Z_{i_1} \dots Z_{i_k} \right). \quad (3.2)$$

What have we accomplished by reducing the problem to a computation in a topological sigma model? Heuristically, the situation is as follows [15]: we have, by reducing to the topological theory, essentially gotten rid of the higher mode excitations of the string. Hence the disk diagram we want is really a sum of Feynman diagrams in a field theory, called holomorphic Chern-Simons theory. Because big Feynman diagrams can be built from smaller ones, we obtain from this way of looking at things combinatorial relations among the correlators, called A_{∞} relations, and it turns out that these determine the correlators uniquely (up to field redefinition). In fact, the A_{∞} relations give a specific algorithm for generating the correlators, and this algorithm generalizes to a more general setting where D -branes are represented as elements of the derived category of coherent sheaves.

We now proceed to flesh out the above heuristic and describe the algorithm in detail. First, we briefly review some mathematical background on A_{∞} products.

3.2 A_{∞} structure

Given a graded vector space B , such as the Dolbeault complex graded by q defined above, an A_{∞} structure on B is defined as a series of products m_k , $k \geq 1$, of degree $2 - k$

$$m_k : B^{\otimes k} \rightarrow B, \tag{3.3}$$

which satisfy the A_∞ constraints:

$$\sum_{r+s+t=n} (-1)^{r+st} m_u(\mathbf{1}^{\otimes r} \otimes m_s \otimes \mathbf{1}^{\otimes t}) = 0, \tag{3.4}$$

for any $n > 0$, where $u = n + 1 - s$. Here we assume the usual sign rule

$$(f \otimes g)(a \otimes b) = (-1)^{|g| \cdot |a|} f(a) \otimes g(b) \tag{3.5}$$

when moving arguments past operators.

The A_∞ products can actually be rephrased in terms of a differential acting on a certain space, with the complicated and unnatural looking relations between them being just the condition that the differential squares to zero [15]. We will not pursue this interpretation here however, except to note that it is useful to consider maps between spaces that commute with the differential. In terms of the A_∞ products, such a map between two spaces B and B' is described as an A_∞ morphism, which is to say it is given by a series of maps

$$f_k : B^{\otimes k} \rightarrow B', \tag{3.6}$$

for $k \geq 1$, which satisfy

$$\sum_{r+s+t=n} (-1)^{r+st} f_u(\mathbf{1}^{\otimes r} \otimes m_s \otimes \mathbf{1}^{\otimes t}) = \sum_{\substack{1 \leq r \leq n \\ i_1 + \dots + i_r = n}} (-1)^q m_r(f_{i_1} \otimes f_{i_2} \otimes \dots \otimes f_{i_r}), \tag{3.7}$$

for any $n > 0$, $u = n + 1 - s$, and $q = (r - 1)(i_1 - 1) + (r - 2)(i_2 - 1) + \dots + (i_{r-1} - 1)$.

Now note that the A_∞ relations give $m_1 \cdot m_1 = 0$, so that B has the structure of a graded differential complex, and we can take cohomology $H^*(B)$. We now come to a theorem that forms the basis for the computational tractability of our results. Let B be as above, except assume that all products m_k are zero for $k \geq 3$ — this structure is called a differential graded algebra (dga). Given an embedding $i : H^*(B) \rightarrow B$ Kadeishvili [34] shows that we may define an A_∞ structure on $H^*(B)$ that has $m_1 = 0$ and an A_∞ morphism f from $H^*(B)$ to B with f_1 equal to the embedding i . Furthermore if B and B' are quasi-isomorphic dga's (that is, there is a map from one to the other that induces an isomorphism on cohomology) then the two Kadeishvili A_∞ structures on $H^*(B)$ and $H^*(B')$ are A_∞ -isomorphic.

There is in fact a well defined algorithm for determining the A_∞ products of Kadeishvili's theorem. The above condition for an A_∞ morphism, for the case $n = 2$, gives

$$im_2 = (i \cdot i) + df_2. \tag{3.8}$$

The cohomology class of the right hand side of the above equation is just that of $i \cdot i$ and hence m_2 is uniquely determined. Therefore df_2 is also uniquely determined, and we can invert d to obtain a (non-unique) choice of f_2 . Now putting $n = 3$ we have

$$im_3 = f_2(\mathbf{1} \otimes m_2) - f_2(m_2 \otimes \mathbf{1}) + (i \cdot f_2) - (f_2 \cdot i) + df_3. \quad (3.9)$$

Once again, this equation uniquely determines m_3 and df_3 , and allows us to make a choice of f_3 . Continuing in this way, it is apparent that all A_∞ products can be determined. The ambiguity in the choice of f_k reflects the ambiguity in the uniqueness clause of the above theorem.

3.3 Holomorphic Chern-Simons Theory

The field theory that the topological B -model on X reduces to is holomorphic Chern-Simons theory:

$$S = \int_X \text{Tr} \left(A \wedge \bar{\partial} A + \frac{2}{3} A \wedge A \wedge A \right) \wedge \Omega, \quad (3.10)$$

where the A is a $(0,1)$ -form on X taking values in $\text{End}(E)$, and Ω is a holomorphic $(3,0)$ -form on X . As mentioned above, computation of the disk correlator in holomorphic Chern-Simons theory reduces to a sum of Feynman diagrams (this reduction can be seen explicitly as localization of the supersymmetric path integral on Feynman fat-graph configurations arising from instantons at infinity, see [35]). The combinatorial relations which the Feynman diagram picture gives rise to are precisely the A_∞ relations. To make a rigorous statement, first define a trace map

$$\gamma(a) = \int_X \text{Tr}(a) \wedge \Omega, \quad (3.11)$$

γ is a degree -3 map in the sense that only $(0,3)$ -forms a have nonzero trace. Define m_1 to be $\bar{\partial}$ and m_2 to be the wedge product together with composition in $\text{End}(E)$ — these give the Dolbeault complex the structure of a dga. The embedding of $\bar{\partial}$ cohomology into the Dolbeault complex by harmonic forms then gives via Kadeishvili an A_∞ structure to $H_{\bar{\partial}}^{0,q}(X, \text{End}(E))$. The correlation functions can then be written [36]:

$$B_{i_0, i_1, \dots, i_k} = \gamma \left(m_2 \left(m_k(\psi_{i_0}, \psi_{i_1}, \dots, \psi_{i_{k-1}}), \psi_{i_k} \right) \right), \quad (3.12)$$

They satisfy the cyclicity property [36]:

$$B_{i_0, i_1, \dots, i_k} = (-1)^{\zeta_k(\zeta_0 + \zeta_1 + \dots + \zeta_{k-1})} B_{i_k, i_0, i_1, \dots, i_{k-1}}. \quad (3.13)$$

which will be important to us later.

Up to now we have been dealing with a single Calabi-Yau filling D -brane. The advantage of working in the above framework is that it extends easily to more general D -brane configurations. For example, (still for a single brane E) we may replace the Dolbeault complex by a Čech complex, thereby turning a difficult problem in analysis, namely inverting $\bar{\partial}$, into a more manageable combinatorial one. The uniqueness theorem above guarantees that the two A_∞ structures obtained are A_∞ -isomorphic. We could also use an injective resolution of a sheaf instead of the Čech complex, and by appropriate abstraction reframe the entire discussion in terms of $\mathbf{D}(X)$. In fact, for now the most convenient complex for

us to use is a hybrid of the Čech complex and that obtained from locally free resolutions (i.e., resolutions by vector bundles). Specifically, we claim that, for a D -brane represented in the derived category by the locally free resolution

$$\mathcal{E}^\bullet = \left(\dots \xrightarrow{d_{n-2}} \mathcal{E}^{n-1} \xrightarrow{d_{n-1}} \mathcal{E}^n \xrightarrow{d_n} \mathcal{E}^{n+1} \xrightarrow{d_{n+1}} \dots \right). \quad (3.14)$$

the following complex has cohomology that gives the correct open string spectrum for the brane and induced A_∞ structure that gives rise to the correct superpotential:

$$\dots \longrightarrow \mathcal{B}^{n-1} \longrightarrow \mathcal{B}^n \longrightarrow \mathcal{B}^{n+1} \longrightarrow \dots, \quad (3.15)$$

where

$$\begin{aligned} \mathcal{B}^n &= \bigoplus_{p+q=n} \mathcal{B}^{p,q} \\ \mathcal{B}^{p,q} &= \check{C}^p(\mathfrak{A}, \mathcal{H}om^q(\mathcal{E}^\bullet, \mathcal{E}^\bullet)). \end{aligned} \quad (3.16)$$

This is shown in [15]. Two points need to be made here. First, the differential in (3.15) is $d = \delta + (-1)^p \mathfrak{d}_q$, with δ the Čech differential and \mathfrak{d}_q given by

$$\mathfrak{d}_n f_{n,p} = \mathfrak{d}_{p+n} \circ f_{n,p} - (-1)^n f_{p+1,n} \circ \mathfrak{d}_p. \quad (3.17)$$

where $\sum_p f_{n,p}$, with $f_{n,p} : \mathcal{E}^p \rightarrow \mathcal{E}^{p+n}$, is an element of

$$\mathcal{H}om^n(\mathcal{E}^\bullet, \mathcal{E}^\bullet) = \bigoplus_p \mathcal{H}om(\mathcal{E}^p, \mathcal{E}^{p+n}). \quad (3.18)$$

Second, to rigorously show that \mathcal{B} indeed reproduces the correct spectrum and superpotential is non-trivial and requires an analysis of elements of the derived category as boundary states of the worldsheet theory [5].

4. Superpotentials for del Pezzo singularities

4.1 Moduli Spaces

Before launching into a more rigorous discussion, let us first consider a heuristic argument that will lead to a conjecture for the form of the superpotential.

First we quickly review the connection of the mathematics of quivers to the physics of D-branes and stability. One should view a quiver as representing a decay of a D-brane. The nodes in the quiver correspond to the decay products, i.e., the so-called “fractional branes” and the arrows correspond to open strings between these decay products. The D-brane we are particularly concerned with is the 3-brane corresponding to a point in X .

At the instant of decay, the open strings corresponding to the arrows should be exactly massless. In the case of B-branes, these masses are a function of the complexified Kähler form $B + iJ$. Here we assume that this masslessness occurs precisely when the del Pezzo surface is collapsed to a point. This assumption was justified in [37].

If one moves away from the critical point where the open strings are massless, then the D-brane may become stable or unstable with respect to the decay. If we deform the Kähler form to some generic value to give a nonzero size the del Pezzo surface (and all the curves within it) then we expect the 3-brane to be stable.

In this resolution one may compute the moduli space of the 3-brane, which should, of course, yield X itself. We need not concern ourselves with the details of this process but we note the following. For more details we refer to [3, 38, 5, 37]. The moduli space of 3-branes is essentially given by the moduli space of representations of the quiver. One takes all possible quiver representations which satisfy “ θ -stability” and then divides by a gauge equivalence.

Physically this moduli space is given by the moduli space of chiral fields (given by the matrices associated to arrows in the quiver) corresponding to classical solutions of the field theory divided by gauge equivalence. Importantly for us, this must mean that *the superpotential imposes conditions on the chiral fields equivalent to the relations in the quiver.*

In other words, finding the critical points of the superpotential must be equivalent to imposing the quiver relations. This leads to an obvious proposition for the superpotential. Let A_i be the chiral fields in the worldvolume gauge theory associated to the arrows in the (non-completed) quiver associated to a del Pezzo surface. The relations will be denoted $r_k(A_1, A_2, \dots) = 0$, where r_k is some polynomial. We know from section 2.3 that each r_k is associated to some arrow in the completed quiver, and so some chiral field R_k . It is believed (see [3, 5], for example) that in terms of the moduli space, setting all R_k equal to zero amounts to restricting the 3-brane to be on the del Pezzo surface S itself. Giving nonzero expectation values to the R_k fields moves the 3-brane off S .

If the superpotential is given by

$$W = \sum_k R_k r_k(A_1, A_2, \dots), \tag{4.1}$$

then, on S , the equations of motion will yield precisely the correct constraints, at least for 3-branes on S . This, therefore, is our conjectured form for the superpotential.

4.2 Quiver Gauge Theories for del Pezzos

Let us now consider more systematically what happens when we put a $D3$ -brane on a shrinking del Pezzo cycle S in a Calabi-Yau X . Now, every BPS space-filling brane corresponds to a topological brane on X , but not vice versa. A point on X is always a valid topological brane; when S is of finite size, the $D3$ -brane will be located on a smooth point of S and as we said we expect it to be BPS. On the other hand, when S shrinks, one can argue (see e.g. [16]) that the $D3$ is marginally stable against decay into the fractional branes introduced earlier. We think of these fractional branes as wrapping S — when S shrinks the point-like $D3$ is allowed to marginally decay to them.

To get a precise description of these fractional branes, we recall that they correspond to the representations L_i , which have resolutions in terms of the projective representations P_i . If we replace the P_i by the corresponding elements of the strongly exceptional collection

(which are all vector bundles in the cases we consider) and use the equivalence between derived categories, we obtain locally free resolutions of the fractional branes. For example in the case of \mathbb{P}^2 and the Beilinson quiver, L_2 is represented in $D(S)$ as

$$\dots \longrightarrow 0 \longrightarrow \mathcal{F}_0^{\oplus 3} \xrightarrow{d-3} \mathcal{F}_1^{\oplus 3} \xrightarrow{d-2} \mathcal{F}_2 \xrightarrow{d-1} 0 \longrightarrow \dots \quad (4.2)$$

The maps in the above complex are determined by the corresponding maps in the quiver resolution of L_2 , using the fact that $\text{Hom}(\mathcal{F}_i, \mathcal{F}_j)$ and $\text{Hom}(P_i, P_j)$ are naturally isomorphic. It turns out that a D3-brane decays into a collection of fractional branes with each fractional brane occurring $\dim \mathcal{F}_i$ times [16]. The quiver gauge theory for a dP_k will thus have $k + 3$ gauge groups, corresponding to each of the L_i , and massless matter in the bi-fundamental from strings stretching between L_i and L_j . The proper setting for discussing the homological structure in this context is an A_∞ category, but we do not need to get so abstract. We simply let \mathcal{M}^\bullet be the direct sum of the locally free resolutions of the L_i , and use it as the starting point for the A_∞ computations. We remind the reader that \mathcal{M}^\bullet is a locally free resolution of sheaves over S , not X .

It will be convenient to represent X as the total space of the normal bundle N of S . Because S is a del Pezzo, $N \cong K_S$ [16]. We are allowed to take this limit in Kähler moduli space because the tree level superpotential does not depend on Kähler parameters.

4.3 From Branes on X to Branes on S

The actual superpotential is computed from the A_∞ products of the Čech complex associated to branes not on S , but on X , i.e., not from $\mathcal{B}^{\bullet,\bullet}$, but rather from the associated complex obtained from considering all the sheaves as embedded in X . The goal of this section is to show that the computation of A_∞ products associated to branes on X essentially reduces to the computation on S . Specifically, we recall that (as we will see in greater detail below) Ext^1 's of sheaves on S considered as sheaves on X include all the Ext^1 's of the sheaves on S plus some extra Ext^1 's, which, after a reversal of arrows, correspond to Ext^2 's of the sheaves on S . We will prove that there is a choice of A_∞ structure over X such that all A_∞ products that contain more than one of the “extra” Ext^1 's vanish. The products that contain no “extra” Ext^1 's are the same as they were over S , and the ones with one “extra” Ext^1 are determined uniquely by the requirement of cyclicity. This, together with the fact that there are no cycles in del Pezzo quivers, will show that the superpotential is linear in the “extra” Ext^1 's, with these Ext^1 's multiplying terms that are just the quiver relations.

To proceed with the proof, let $\pi : E \rightarrow S$ be the projection from the total space of K_S to the del Pezzo S . We have a canonical section $\mathcal{O} \rightarrow \pi^*K_S$, given as follows: to each point in E we tautologically associate a point of S and an element of the fiber of K_S over that point; this element can be viewed as an element of the fiber of π^*K_S over the original point in E . Dualizing, we get a canonical map which fits into an exact sequence of sheaves

$$0 \rightarrow \pi^*(K_S^*) \rightarrow \mathcal{O} \rightarrow \mathcal{O}_S \rightarrow 0. \quad (4.3)$$

and we can think of the first two terms as a locally free resolution of \mathcal{O}_S . The key point now is to tensor the resolution \mathcal{M}^\bullet with the above resolution of \mathcal{O}_S in order to obtain a locally free resolution of i_*M :

$$\begin{array}{ccccccc}
 & \longrightarrow & \pi^{-1}\mathcal{M}^0 & \longrightarrow & \pi^{-1}\mathcal{M}^1 & \longrightarrow & \pi^{-1}\mathcal{M}^2 \longrightarrow \\
 \uparrow & & \uparrow d^1 & & \uparrow & & \uparrow \\
 j & & & & & & \\
 & \longrightarrow & \pi^*(K_S^*) \otimes \pi^{-1}\mathcal{M}^0 & \xrightarrow{d^0} & \pi^*(K_S^*) \otimes \pi^{-1}\mathcal{M}^1 & \longrightarrow & \pi^*(K_S^*) \otimes \pi^{-1}\mathcal{M}^2 \longrightarrow \\
 & & & & i & & \\
 & \longrightarrow & & & & &
 \end{array} \tag{4.4}$$

Collapsing the above double complex along the diagonal we get a free resolution, and the associated spectral sequence, which collapses at the E_2 term, shows that it is in fact a resolution of $i_*\mathcal{M}$. This can also be viewed as the Cone construction [5]. We will choose to retain the bi-grading, so let us represent the above resolution as $\mathcal{M}^{\bullet,\bullet}$ (the first index corresponds to the index of \mathcal{M}^\bullet , and the second is either 0, for $\pi^{-1}\mathcal{M}^\bullet$, or 1, for $\pi^*(K_S^*) \otimes \pi^{-1}\mathcal{M}^\bullet$). We now define

$$\mathcal{C}^{p,i,j} = \check{C}^p(\pi^{-1}\mathcal{U}, \mathcal{H}om^{i,j}(\mathcal{M}^{\bullet,\bullet}, \mathcal{M}^{\bullet,\bullet})) \tag{4.5}$$

Here \mathcal{U} is an affine open cover of S , and hence $\pi^{-1}\mathcal{U}$ is an affine open cover of E . The complex \mathcal{C} is central in our analysis. There are several differentials we define on \mathcal{C} . First of all, in the locally free resolution $\mathcal{M}^{\bullet,\bullet}$ label the differentials that increase the first index by $d^0_{i,j}$ and that which increases the second index by $d^1_{i,j}$. The combination $d_{i,j} = d^0_{i,j} + (-1)^i d^1_{i,j}$ is the standard differential associated to the locally free resolution of \mathcal{M} over E . Now, given a section of $\mathcal{H}om^{i,j}(\mathcal{M}^\bullet, \mathcal{M}^\bullet)$ over an open set U , i.e., a section of $\bigoplus_{p,q} \text{Hom}_U(\mathcal{M}^{p,q}, \mathcal{M}^{p+i,q+j})$, we can denote it by

$$\sum_{p,q} f_{p,q}^{i,j} \tag{4.6}$$

where

$$f_{p,q}^{i,j} : \mathcal{M}^{p,q}(U) \rightarrow \mathcal{M}^{p+i,q+j}(U) \tag{4.7}$$

Then we can define differentials

$$\mathfrak{d}^0_{i,j} f_{p,q}^{i,j} = d^0_{i,j} f_{p,q}^{i,j} - (-1)^{(i+j)} f_{p+1,q}^{i,j} d^0_{i,j} \tag{4.8}$$

$$\mathfrak{d}^1_{i,j} f_{p,q}^{i,j} = d^1_{i,j} f_{p,q}^{i,j} - (-1)^{(i+j)} f_{p,q+1}^{i,j} d^1_{i,j}. \tag{4.9}$$

Finally, we have the Čech differential δ , and we define the total differential on \mathcal{C} by $d = \delta + (-1)^p (\mathfrak{d}^0 + (-1)^i \mathfrak{d}^1)$. Now, the sum $\mathfrak{d}^0 + (-1)^i \mathfrak{d}^1$ is the standard differential associated to the locally free resolution of \mathcal{M} over E , so that collapsing on the (i, j) indices yields the double complex in [15], showing that \mathcal{C} does indeed correctly compute the A_∞ products.

To actually get a handle on determining the A_∞ algebra, it is useful to collapse \mathcal{C} in a different way and leverage our knowledge of the A_∞ structure for sheaves on the del Pezzo S . Specifically, let us collapse the complex on the (p, i) indices:

$$\mathcal{D}^{q,j} = \bigoplus_{p+i=q} \mathcal{C}^{p,i,j} \tag{4.10}$$

$\mathcal{D}^{\bullet,\bullet}$ is a double complex with anticommuting differentials $d_0 = \delta + (-1)^p \mathfrak{d}^0$, which increases the first index, and $d_1 = (-1)^q \mathfrak{d}^1$, which increases the second one, that add up to d . The desired cohomology is computed using a spectral sequence associated to this double complex, which by arguments of [16] degenerates at the E_2 term to give:

$$\begin{array}{cccc}
 & 0 & & 0 & & 0 & & 0 \\
 & & \uparrow j & & & & & \\
 \text{Ext}_S^{-1}(M, M \otimes K_S) & & \text{Ext}_S^0(M, M \otimes K_S) & & \text{Ext}_S^1(M, M \otimes K_S) & & \text{Ext}_S^2(M, M \otimes K_S) & \\
 \text{Ext}_S^{-1}(M, M) & & \text{Ext}_S^0(M, M) & & \text{Ext}_S^1(M, M) & & \text{Ext}_S^2(M, M) & \\
 \hline
 & & & q & & & &
 \end{array} \tag{4.11}$$

(In our exposition the E_1 term is given by taking cohomology with respect to d_1). Serre duality shows that $\text{Ext}_S^i(M, M) \cong \text{Ext}_S^{2-i}(M, M \otimes K_S)$, so that

$$\text{Ext}_X^1(i_*M, i_*M) \cong \text{Ext}_S^1(M, M) \oplus \text{Ext}_S^2(M, M). \tag{4.12}$$

The two terms on the right correspond to the Ext^1 s and “extra” Ext^1 s, respectively.

Now, the bottom row in the above diagram reproduces the cohomology of the complex $\mathcal{B}^{p,i}$ — the complex associated to branes on S rather than X . In fact, we may naturally embed $\mathcal{B}^{p,i}$ in

$$\check{C}^p(\pi^{-1}\mathcal{U}, \mathcal{H}om^i(\pi^{-1}\mathcal{M}^\bullet, \pi^{-1}\mathcal{M}^\bullet)). \tag{4.13}$$

This complex, in turn, can be viewed as a sub-complex of $\mathcal{C}^{p,i,0}$. To see why, note that, from (4.4), a section of $\mathcal{H}om^i(\pi^{-1}\mathcal{M}^\bullet, \pi^{-1}\mathcal{M}^\bullet)$ determines a section of $\mathcal{H}om^{i,0}(\mathcal{M}^{\bullet,\bullet}, \mathcal{M}^{\bullet,\bullet})$; basically it gives directly the maps among the sheaves in the upper row of (4.4), and, taking the identity map on $\pi^*(K_S^*) \otimes \pi^{-1}\mathcal{M}^\bullet$, it determines the maps for the sheaves in the lower row as well. Also, from (4.11), we see that the composition of these embeddings induces an isomorphism from the cohomology of $\mathcal{B}^{p,i}$ to the $j = 0$ part of the cohomology of $\mathcal{C}^{p,i,j}$.

Now, in order to carry out the A_∞ procedure we must choose representatives for all cohomology classes in $\mathcal{C}^{p,i,j}$. The upshot of the construction in the previous paragraph is that it gives us a natural choice of representatives of the $j = 0$ part of the cohomology; in fact, it shows that the A_∞ products of these $j = 0$ cohomology classes are exactly the same as those in $\mathcal{B}^{p,i}$. In other words, for the $j = 0$ generators the A_∞ products are just those defined over S . This accomplishes part of our goal of reducing the computation over X to a computation over S ; to finish we have to deal with products that may contain some $j = 1$ generators.

The $j = 1$ cohomology generators are the ones that contribute the “extra” Ext^1 s. To carry out the A_∞ procedure, we must pick representatives of their cohomology classes. We choose these to be homogeneous of j degree 1, or, in other words, to lie in

$$\check{C}^p(\pi^{-1}\mathfrak{U}, \mathcal{H}om^i(\pi^{-1}\mathcal{M}^\bullet \otimes \pi^{-1}K_S, \pi^{-1}\mathcal{M}^\bullet)) \quad (4.14)$$

Clearly this is the most natural choice, though it should be pointed out that we could have done something stupid and chosen the generator to have a nonzero (exact) $j = 0$ part, for example. The advantage of having homogeneous generators is that their products are homogeneous as well, and so vanish if they have $j > 1$.

Now we claim that any m_k that contains more than one $j = 1$ generator vanishes. The naive argument would invoke the j grading and the fact that there are no elements in $\mathcal{E}^{p,i,j}$ with $j > 1$. The obvious flaw is the fact that m_k does not respect the overall grading — it in fact has degree $2 - k$. Thus we have to be more careful. We claim that although m_k has nonzero degree with respect to the overall grading $p + i + j$, through a careful choice of f_k , which we now construct, we can make m_k respect the j grading. The claim at the top of this paragraph then immediately follows.

We show by induction that all the f_k and m_k respect the j grading. Clearly f_1 and m_1 respect the j grading. Suppose that this is also true for all $k \leq n$. We have for all $n + 1$ (3.7), which can be rewritten as an equation determining m_{n+1} in terms of the lower m_k and f_k (for $n + 1 = 3$, for example, this is (3.9)). So we immediately see that it’s true for m_{n+1} . We now deal with f_{n+1} . We suppress its arguments, but everywhere below f_{n+1} and df_{n+1} will stand for f_{n+1} and df_{n+1} applied to their arguments. Now equation (3.7) again gives df_{n+1} as an expression in terms of m_{n+1} and the m_k and f_k for $k \leq n$. We have to make a choice of f_{n+1} that respects the j grading, i.e., we want f_{n+1} to have the same j degree as df_{n+1} . Now, the case when all the arguments have $j = 0$ has been discussed above and clearly we have already defined f_{n+1} to have $j = 0$. When more than one argument has $j = 1$ then, because all the terms in the expression for df_{n+1} are homogeneous, we have $df_{n+1} = 0$, so that we can choose $f_{n+1} = 0$. The nontrivial case is when exactly one argument has $j = 1$. In that case, df_{n+1} is homogeneous of j degree 1 and hence lies in the sub-complex

$$\mathcal{E}^{p,i,1} = \check{C}^p(\pi^{-1}\mathfrak{U}, \mathcal{H}om^i(\pi^{-1}\mathcal{M}^\bullet \otimes \pi^{-1}K_S, \pi^{-1}\mathcal{M}^\bullet)). \quad (4.15)$$

The crucial point is now that the embedding $\mathcal{E}^{p,i,1} \subset \mathcal{E}^{p,i,j}$ induces an injection in cohomology. This can easily be seen from the spectral sequence (4.11) — the cohomology of $\mathcal{E}^{p,i,1}$ reproduces precisely the upper, $j = 1$ row in the diagram. Therefore df_{n+1} is exact not only in $\mathcal{E}^{p,i,j}$ but also in the sub-complex $\mathcal{E}^{p,i,1}$. Therefore we can choose f_{n+1} to be in $\mathcal{E}^{p,i,1}$, so that it will have $j = 1$. Thus we see that we can always choose f_{n+1} to respect the j grading. This completes the inductive step.

Together with the cyclicity property this determines all the A_∞ products in $\mathcal{E}^{p,i,j}$ in terms of those over S . To restate, we have the original A_∞ algebra reproduced when all the arguments are the original Ext^1 s (i.e., have $j = 0$), any product that involves more than one “extra” Ext^1 (i.e., one that has $j = 1$) must vanish, while any product that contains

exactly one “extra” Ext^1 is determined uniquely by requiring it to reproduce correlators that obey the cyclicity property (3.13). Having accomplished the reduction and thus shown that the superpotential is linear in the “extra” Ext^1 s, we now determine the A_∞ products over S and relate them to the quiver relations.

5. A_∞ relations and quivers

We must determine the A_∞ products over S , i.e., those of $\mathcal{B}^{\bullet,\bullet}$, defined in (3.16). We know by Bondal’s theorem that $\mathbf{D}(S)$ is equivalent to $\mathbf{D}(A\text{-mod})$, where A is the path algebra of the associated quiver. The operational version of this equivalence that will suffice for us is as follows. First, construct a complex of quiver representations M^\bullet by summing the projective resolutions of the L_i . In the usual way it gives rise to the graded dga $\text{End}(M)^\bullet$. There is a natural map of this complex into \mathcal{B} given by interpreting maps in $\text{End}(M)$ as global sections of the $\text{Hom}(\mathcal{F}_i, \mathcal{F}_j)$ and mapping them to Čech 0-cochains. Because these 0-cochains are global sections, they are annihilated by the Čech part of the differential, and one thus quickly sees that this map is a map of dga’s. In fact, (the derivation of) Bondal’s theorem shows that it is a quasi-isomorphism. This allows us to apply Kadeishvili’s theorem and compute the A_∞ structure of \mathcal{B} in the quiver dga $\text{End}(M)^\bullet$.

We will see that we obtain a form of the superpotential exactly as conjectured in section 4.1.

5.1 A simple example

We start with the Beilinson quiver, corresponding to $S = \mathbb{P}^2$:

$$\begin{array}{ccccc}
 & & a_0 & & b_0 \\
 & \swarrow & \leftarrow & \searrow & \\
 \circ & & & & \circ \\
 & \nwarrow & a_1 & \nearrow & \\
 & & & & \\
 & \swarrow & a_2 & \searrow & \\
 \circ & & & & \circ \\
 & \nwarrow & & \nearrow & \\
 & & v_1 & & v_2
 \end{array} \tag{5.1}$$

with relations $a_\alpha b_\beta = a_\beta b_\alpha$. We recall that we have the projective resolutions:

$$\begin{aligned}
 0 &\longrightarrow P_0 \longrightarrow L_0 \longrightarrow 0 & (5.2) \\
 0 &\longrightarrow P_0^{\oplus 3} \longrightarrow P_1 \longrightarrow L_1 \longrightarrow 0 \\
 0 &\longrightarrow P_0^{\oplus 3} \longrightarrow P_1^{\oplus 3} \longrightarrow P_2 \longrightarrow L_2 \longrightarrow 0
 \end{aligned}$$

and that

$$\begin{aligned}
 \dim \text{Ext}^1(L_i, L_j) &= n_{ij} \\
 \dim \text{Ext}^2(L_i, L_j) &= r_{ij},
 \end{aligned} \tag{5.3}$$

where n_{ij} counts arrows and r_{ij} counts relations.

We start by choosing specific generators of the Ext^i . Recalling that the Ext^\bullet can be represented as morphisms between resolutions of the L_i , we can choose the three generators a_i of $\text{Ext}^1(L_1, L_0)$ to be

$$\begin{array}{ccc}
 P_0^{\oplus 3} & \longrightarrow & P_1 \\
 \downarrow \pi_i & & \\
 P_0 & &
 \end{array} \tag{5.4}$$

where π_i is projection on the i 'th factor. As far as the generators \mathbf{b}_i of $\text{Ext}^1(L_2, L_1)$ we have \mathbf{b}_0 represented as

$$\begin{array}{ccccc} P_0^{\oplus 3} & \longrightarrow & P_1^{\oplus 3} & \longrightarrow & P_2 \\ \downarrow \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -1 \\ 0 & 1 & 0 \end{pmatrix} & & \downarrow (1 \ 0 \ 0) & & \\ P_0^{\oplus 3} & \longrightarrow & P_1 & & \end{array} \quad (5.5)$$

and the other \mathbf{b}_i represented similarly. We also have the relations r_i in $\text{Ext}^2(L_3, L_2)$, represented by

$$\begin{array}{ccccc} P_0^{\oplus 3} & \longrightarrow & P_1^{\oplus 3} & \longrightarrow & P_2 \\ \downarrow \pi_i & & & & \\ P_0 & & & & \end{array} \quad (5.6)$$

Clearly, the only potentially nontrivial products are $m_2(\mathbf{a}_j, \mathbf{b}_i)$, and one easily sees by composing the representatives for \mathbf{a}_j and \mathbf{b}_i that $m_2(\mathbf{a}_j, \mathbf{b}_i) = \epsilon^{ijk} r_k$. This gives rise to the superpotential

$$W = \epsilon^{ijk} A_i B_j R_k \quad (5.7)$$

which is the correct superpotential for this quiver gauge theory on the orbifold $\mathbb{C}^3/\mathbb{Z}_3$ [1]. Again, the A_i and B_j are massless moduli corresponding to the internal structure of the shrinking cycle \mathbb{P}^2 , while R_k is the modulus that corresponds to moving the $D3$ -brane off the singularity. We note that the superpotential is of the desired form, linear in the ‘‘extra’’ Ext^1 s R_i , which multiply the relations. To write it out explicitly, we have

$$W = (A_0 B_1 - B_1 A_0) R_2 + (A_1 B_2 - B_2 A_1) R_0 + (A_2 B_0 - B_0 A_2) R_1 \quad (5.8)$$

5.2 del Pezzo 1

Let us consider the quiver associated to $d\mathbb{P}_1$. It is:

$$\quad (5.9)$$

subject to the relations $r_0 = b_0 d_1 - b_1 d_0$, $s_0 = a b_0 d_2 - c d_0 = 0$, and $s_1 = a b_1 d_2 - c d_1 = 0$. Denote the corresponding generators of Ext^2 by r_0 , s_0 , and s_1 . We first pick maps of projective resolutions representing these generators, which all turn out to be uniquely determined. We have the projective resolutions:

$$\begin{array}{ccccccc} 0 & \longrightarrow & P_0 & \longrightarrow & L_0 & \longrightarrow & 0 \\ 0 & \longrightarrow & P_0 & \longrightarrow & P_1 & \longrightarrow & L_1 \longrightarrow 0 \\ 0 & \longrightarrow & P_0 \oplus P_1^{\oplus 2} & \longrightarrow & P_2 & \longrightarrow & L_2 \longrightarrow 0 \\ 0 & \longrightarrow & P_0^{\oplus 2} \oplus P_1 & \longrightarrow & P_2^{\oplus 3} & \longrightarrow & P_3 \longrightarrow L_3 \longrightarrow 0 \end{array} \quad (5.10)$$

We choose representatives of $\text{Ext}^1(L_i, L_j)$ as follows: for $i \leq 3$, there are no relations originating at the i 'th node of the quiver and hence the maps representing \mathbf{b}_i , \mathbf{a} , and \mathbf{c} are uniquely determined. The choice of representative of \mathbf{d}_i is uniquely determined as well. That is to say, in the diagram

$$\begin{array}{ccccc}
 P_0^{\oplus 2} \oplus P_1 & \longrightarrow & P_2^{\oplus 3} & \longrightarrow & P_3 \\
 \downarrow & & \downarrow & & \\
 P_0 \oplus P_1^{\oplus 2} & \longrightarrow & P_2 & &
 \end{array} \tag{5.11}$$

the bottom horizontal map is injective, so that the left vertical map is uniquely determined by the right vertical map, which we take to be projection on the i 'th factor. Finally, for the generators of Ext^2 we take the obvious uniquely determined maps from the projective resolution of L_3 to the other L_i .

We want to compute all products m_k of the various Ext^1 's. Such products will all be in Ext^2 , and because only $\text{Ext}^2(L_3, L_0)$ and $\text{Ext}^2(L_3, L_1)$ are nonzero we see that the possible nonzero products are $m_2(\mathbf{b}_i, \mathbf{d}_j)$, $m_2(\mathbf{c}, \mathbf{d}_i)$, and $m_3(\mathbf{a}, \mathbf{b}_i, \mathbf{d}_j)$. Let's look at the first of these; the relevant composition is:

$$\begin{array}{ccccc}
 P_0^{\oplus 2} \oplus P_1 & \longrightarrow & P_2^{\oplus 3} & \longrightarrow & P_3 \\
 \downarrow & & \downarrow & & \\
 P_0 \oplus P_1^{\oplus 2} & \longrightarrow & P_2 & & \\
 \downarrow & & & & \\
 P_0 & \longrightarrow & P_1 & &
 \end{array} \tag{5.12}$$

The map from the first to the second row is \mathbf{d}_j and that from the second to the third row is \mathbf{b}_i . For convenience, we label the individual P_k s that occur in various parts of the diagram. The upper left entry is $P_0^{\oplus 2} \oplus P_1$, where each summand corresponds to a different relation. We naturally label the two P_0 s as S_0 and S_1 , and we label the P_1 as R_0 . The entry below this one is $P_0 \oplus P_1^{\oplus 2}$, and here each P_k corresponds to an arrow emanating from P_2 . Thus we label the P_0 as C and the two P_1 s as B_0 and B_1 .

Let us consider the possible maps we can have. For $k > l$ there are no nonzero maps from P_k to P_l . From each P_k to itself there is the identity map, and it is the only one that will be of use to us. For $k < l$, however, there are several ways to map P_k to P_l , each corresponding to a path from node l to node k . Thus, for example, there is one map from P_0 to P_1 , denoted by a .

To see how \mathbf{d}_j acts note that it maps $R_0 \oplus S_0 \oplus S_1$ to $C \oplus B_0 \oplus B_1$. We can thus represent its action on $R_0 \oplus S_0 \oplus S_1$ as a 3 by 3 matrix. From the definition of \mathbf{d}_j it is easy to obtain (by slight abuse of notation we denote by \mathbf{d}_j both itself and its restriction to $R_0 \oplus S_0 \oplus S_1$):

$$\mathbf{d}_0 = \begin{pmatrix} 0 & -1 & 0 \\ 0 & 0 & 0 \\ -1 & 0 & 0 \end{pmatrix}, \mathbf{d}_1 = \begin{pmatrix} 0 & 0 & -1 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}, \mathbf{d}_2 = \begin{pmatrix} 0 & 0 & 0 \\ 0 & a & 0 \\ 0 & 0 & a \end{pmatrix}. \quad (5.13)$$

Note that these matrix elements are simply obtained by “contracting” the relevant relation (which indexes the column) with d_j to obtain the elements of the column. Similar reasoning shows that the \mathbf{b}_i act as follows:

$$\mathbf{b}_0 = (0 \ 1 \ 0), \mathbf{b}_1 = (0 \ 0 \ 1). \quad (5.14)$$

The nonzero compositions are

$$\mathbf{b}_0\mathbf{d}_1 = (1 \ 0 \ 0), \mathbf{b}_0\mathbf{d}_2 = (0 \ a \ 0), \quad (5.15)$$

$$\mathbf{b}_1\mathbf{d}_0 = (-1 \ 0 \ 0), \mathbf{b}_1\mathbf{d}_2 = (0 \ 0 \ a), \quad (5.16)$$

Now, note that the compositions $\mathbf{b}_0\mathbf{d}_2$ and $\mathbf{b}_1\mathbf{d}_2$ can both be factored through the leftmost map $P_0 \rightarrow P_1$, and hence are exact in the quiver dga. So the only nonzero products are

$$m_2(\mathbf{b}_0, \mathbf{d}_1) = r_0, m_2(\mathbf{b}_1, \mathbf{d}_0) = -r_0. \quad (5.17)$$

A good shorthand way of expressing this result is that $m_2(\mathbf{b}_i, \mathbf{d}_j) = (b_i d_j, R_0)r_0 + (b_i d_j, S_0)s_0 + (b_i d_j, S_1)s_1$, where the parentheses denote the coefficient of the string represented by the left argument in the relation represented by the right argument. As one traces through the above manipulations it is clear that this is a general result that always holds when one computes the products m_2 . Thus we have $m_2(\mathbf{c}, \mathbf{d}_0) = -s_0$, $m_2(\mathbf{c}, \mathbf{d}_1) = -s_1$, and $m_2(\mathbf{c}, \mathbf{d}_2) = 0$.

To compute m_3 , we first have to define a choice of f_2 , which must satisfy

$$im_2 = (i \cdot i) + df_2. \quad (5.18)$$

Quick inspection shows that we may take $f_2 = 0$ everywhere except for $f_2(\mathbf{b}_i, \mathbf{d}_j)$. From the above analysis, we see that

$$df_2(\mathbf{b}_0, \mathbf{d}_2) = (0 \ -a \ 0), df_2(\mathbf{b}_1, \mathbf{d}_2) = (0 \ 0 \ -a). \quad (5.19)$$

where a denotes right multiplication and the notation indicates a map $P_0^{\oplus 2} \oplus P_1 \rightarrow P_1$. Hence we can define $f_2(\mathbf{b}_0, \mathbf{d}_2)$ and $f_2(\mathbf{b}_1, \mathbf{d}_2)$ respectively as:

$$\begin{array}{ccccc} P_0^{\oplus 2} \oplus P_1 & \longrightarrow & P_2^{\oplus 3} & \longrightarrow & P_3 \\ \downarrow (0 \ -1 \ 0) & & \downarrow 0 & & \\ P_0 & \longrightarrow & P_1 & & \end{array} \quad (5.20)$$

$$\begin{array}{ccc}
 P_0^{\oplus 2} \oplus P_1 & \longrightarrow & P_2^{\oplus 3} \longrightarrow P_3 \\
 \downarrow (0 \ 0 \ -1) & & \downarrow 0 \\
 P_0 & \longrightarrow & P_1
 \end{array} \tag{5.21}$$

Now, we have

$$im_3 = f_2(\mathbf{1} \otimes m_2) - f_2(m_2 \otimes \mathbf{1}) + (i \cdot f_2) - (f_2 \cdot i) + df_3, \tag{5.22}$$

so that, recalling the sign rule (3.5), $m_3(\mathbf{a}, \mathbf{b}_i, \mathbf{d}_j) = -[\mathbf{a} \cdot f_2(\mathbf{b}_i, \mathbf{d}_j)]$. Composing with \mathbf{a} we see that $m_3(\mathbf{a}, \mathbf{b}_0, \mathbf{d}_2) = \mathbf{s}_0$ and $m_3(\mathbf{a}, \mathbf{b}_1, \mathbf{d}_2) = \mathbf{s}_1$. Again, a shorthand way of expressing this result is $m_3(\mathbf{a}, \mathbf{b}_i, \mathbf{d}_j) = (ab_i d_j, R_0)r_0 + (ab_i d_j, S_0)s_0 + (ab_i d_j, S_1)s_1$. These are all the nonzero A_∞ products in this example.

One can see by carrying out the A_∞ algorithm that this formula generalizes to all the m_k in any quiver. We sketch the argument modulo various signs, which have to be checked carefully. Let the relations be labeled by R_i and the corresponding generators of Ext^2 by r_i , as above. One proceeds by induction on k_0 . Let's take the following inductive hypothesis: for all $j < k_0$, we have

$$m_j(a_1, \dots, a_j) = \sum_i (a_1 \dots a_j, R_i)r_i \tag{5.23}$$

as well as a statement about f_k for which we need to introduce some notation. Each a_i is an Ext^1 , so that it can be represented as a map between the projective resolution of $L_{m(i)}$ and $L_{m(i-1)}$. According to this notation, a_i is an arrow between node $m(i)$ and $m(i-1)$. The projective resolution of $L_{m(i)}$ is

$$\dots \longrightarrow \bigoplus_k P_k^{\oplus r_{m(i)k}} \longrightarrow \bigoplus_k P_k^{\oplus n_{m(i)k}} \longrightarrow P_{m(i)} \longrightarrow L_{m(i)} \longrightarrow 0. \tag{5.24}$$

The second part of the inductive hypothesis is that for $j < k_0$, $f_j(a_1, \dots, a_j)$ is represented as:

$$\begin{array}{ccc}
 \bigoplus_k P_k^{\oplus r_{m(j)k}} & \longrightarrow & \bigoplus_k P_k^{\oplus n_{m(j)k}} \longrightarrow P_{m(j)} \\
 \downarrow & & \downarrow 0 \\
 \bigoplus_k P_k^{\oplus n_{m(0)k}} & \longrightarrow & P_{m(0)}
 \end{array} \tag{5.25}$$

where the left vertical map takes each relation to an Ext^1 determined by the contraction of the relation with $a_1 \dots a_j$ (we extract only the linear terms in the contraction, as only these correspond to Ext^1 s). f_j with any Ext^2 s as arguments vanish.

For the inductive step, we have to prove the analogous statements for m_{k_0} . We use (3.7) to write m_{k_0} in terms of the lower m_j s and f_j s. We then note from the above form of the f_j that all terms vanish except $(i \cdot f_{k_0-1})$ — basically because the only nonzero map in (5.25) takes relations to Ext^1 s, so the composition of two f_j s vanishes. This straightforwardly leads to (5.23) for m_{k_0} . Inverting d shows that f_{k_0} may be chosen to be of the form (5.25).

Thus essentially one knows these products as soon as one knows all the relations in the quiver. It follows that the term in the superpotential that multiplies the “extra” Ext^1 corresponding to a given relation is simply that relation (written as a polynomial in the Ext^1 s). We can now write down the superpotential:

$$W = R_0(B_0D_1 - B_1D_0) + S_0(AB_0D_2 - CD_0) + S_1(AB_1D_2 - CD_1) \quad (5.26)$$

6. Conclusions

We have given an effective method for computing superpotentials for quiver gauge theories associated with shrinking del Pezzo cycles. We showed that the superpotential is linear in the fields that correspond to Ext^2 s in the del Pezzo quiver, and that each such field multiplies a polynomial which is just the corresponding relation. To do this we performed a precise reduction of the problem from one involving sheaves on the Calabi-Yau to one involving sheaves on the del Pezzo. We solved the problem on the del Pezzo by switching to the algebraically more tractable category of quiver representations and explicitly evaluating the A_∞ products there. We did only the cases where S is \mathbb{P}^2 and dP_1 , but these examples show that the algorithm is trivial to carry out provided one has the quiver and the relations. These, of course, might not be so trivial to obtain, especially for the higher del Pezzos, which themselves have complex structure moduli. These complex structure moduli are contained in the choice of points to be blown up on \mathbb{P}^2 , and will show up in the quiver relations.

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References

- [1] M.R. Douglas and G.W. Moore, *D-branes, quivers and ALE instantons*, hep-th/9603167.
- [2] D.-E. Diaconescu, M.R. Douglas and J. Gomis, *Fractional branes and wrapped branes*, *JHEP* **02** (1998) 013 [hep-th/9712230].
- [3] M.R. Douglas, B. Fiol and C. Romelsberger, *The spectrum of BPS branes on a noncompact Calabi-Yau*, *JHEP* **09** (2005) 057 [hep-th/0003263].
- [4] M.R. Douglas, *Lectures on D-branes on Calabi-Yau manifolds*, Prepared for ICTP Spring School on Superstrings and Related Matters, Trieste, Italy, 2-10 Apr 2001, http://www.ictp.trieste.it/~pub_off/lectures/lms007/Douglas/Douglas.pdf.
- [5] P.S. Aspinwall, *D-branes on Calabi-Yau manifolds*, to appear in the proceedings of TASI 2003, hep-th/0403166.
- [6] S. Franco, A. Hanany, K.D. Kennaway, D. Vegh and B. Wecht, *Brane dimers and quiver gauge theories*, *JHEP* **01** (2006) 096 [hep-th/0504110].

- [7] C.V. Johnson and R.C. Myers, *Aspects of type-IIB theory on ALE spaces*, *Phys. Rev. D* **55** (1997) 6382 [[hep-th/9610140](#)].
- [8] I.R. Klebanov and E. Witten, *Superconformal field theory on threebranes at a Calabi-Yau singularity*, *Nucl. Phys. B* **536** (1998) 199 [[hep-th/9807080](#)].
- [9] D.R. Morrison and M.R. Plesser, *Non-spherical horizons, I*, *Adv. Theor. Math. Phys.* **3** (1999) 1 [[hep-th/9810201](#)].
- [10] B. Feng, A. Hanany and Y.-H. He, *D-brane gauge theories from toric singularities and toric duality*, *Nucl. Phys. B* **595** (2001) 165 [[hep-th/0003085](#)].
- [11] B. Feng, A. Hanany and Y.-H. He, *Phase structure of D-brane gauge theories and toric duality*, *JHEP* **08** (2001) 040 [[hep-th/0104259](#)].
- [12] B. Feng, A. Hanany, Y.-H. He and A.M. Uranga, *Toric duality as Seiberg duality and brane diamonds*, *JHEP* **12** (2001) 035 [[hep-th/0109063](#)].
- [13] B. Feng, S. Franco, A. Hanany and Y.-H. He, *Unhiggsing the Del Pezzo*, *JHEP* **08** (2003) 058 [[hep-th/0209228](#)].
- [14] M. Wijnholt, *Large volume perspective on branes at singularities*, *Adv. Theor. Math. Phys.* **7** (2004) 1117 [[hep-th/0212021](#)].
- [15] P.S. Aspinwall and S. Katz, *Computation of superpotentials for D-branes*, *Commun. Math. Phys.* **264** (2006) 227 [[hep-th/0412209](#)].
- [16] P.S. Aspinwall and I.V. Melnikov, *D-branes on vanishing Del Pezzo surfaces*, *JHEP* **12** (2004) 042 [[hep-th/0405134](#)].
- [17] C. Beasley, B.R. Greene, C.I. Lazaroiu and M.R. Plesser, *D3-branes on partial resolutions of abelian quotient singularities of Calabi-Yau threefolds*, *Nucl. Phys. B* **566** (2000) 599 [[hep-th/9907186](#)].
- [18] C.E. Beasley and M.R. Plesser, *Toric duality is Seiberg duality*, *JHEP* **12** (2001) 001 [[hep-th/0109053](#)].
- [19] B. Feng, A. Hanany, Y.H. He and A. Iqbal, *Quiver theories, soliton spectra and Picard-Lefschetz transformations*, *JHEP* **02** (2003) 056 [[hep-th/0206152](#)].
- [20] A. Hanany and A. Iqbal, *Quiver theories from D6-branes via mirror symmetry*, *JHEP* **04** (2002) 009 [[hep-th/0108137](#)].
- [21] F. Cachazo, B. Fiol, K.A. Intriligator, S. Katz and C. Vafa, *A geometric unification of dualities*, *Nucl. Phys. B* **628** (2002) 3 [[hep-th/0110028](#)].
- [22] C.P. Herzog and J. Walcher, *Dibaryons from exceptional collections*, *JHEP* **09** (2003) 060 [[hep-th/0306298](#)].
- [23] C.P. Herzog, *Exceptional collections and Del Pezzo gauge theories*, *JHEP* **04** (2004) 069 [[hep-th/0310262](#)].
- [24] M.R. Douglas, *D-branes, categories and $N = 1$ supersymmetry*, *J. Math. Phys.* **42** (2001) 2818 [[hep-th/0011017](#)].
- [25] P.S. Aspinwall and A.E. Lawrence, *Derived categories and zero-brane stability*, *JHEP* **08** (2001) 004 [[hep-th/0104147](#)].

- [26] A.I. Bondal, *Representations of associative algebras and coherent sheaves*, *Math. USSR Izvestiya* **34** (1990) 23–42.
- [27] A.I. Bondal, *Helices, representations of quivers and Koszul algebras*, in A.N. Rudakov et al., eds., *Helices and vector bundles*, LMS Lecture Notes **148**, pages 75–96, Cambridge, 1990.
- [28] C.P. Herzog, *Seiberg duality is an exceptional mutation*, *JHEP* **08** (2004) 064 [[hep-th/0405118](#)].
- [29] R. Hartshorne, *Algebraic geometry*, *Graduate Texts in Mathematics* **52**, Springer-Verlag, 1977.
- [30] P. Seidel and R.P. Thomas, *Braid groups actions on derived categories of coherent sheaves*, *Duke Math. J.* **108** (2001) 37–108, [math.AG/0001043](#).
- [31] S. Katz and E. Sharpe, *D-branes, open string vertex operators and ext groups*, *Adv. Theor. Math. Phys.* **6** (2003) 979 [[hep-th/0208104](#)].
- [32] P.S. Aspinwall, *A point's point of view of stringy geometry*, *JHEP* **01** (2003) 002 [[hep-th/0203111](#)].
- [33] I. Brunner, M.R. Douglas, A.E. Lawrence and C. Römelsberger, *D-branes on the quintic*, *JHEP* **08** (2000) 015 [[hep-th/9906200](#)].
- [34] T.V. Kadeishvili, *The algebraic structure in the homology of an A_∞ -algebra*, *Soobshch. Akad. Nauk. Gruzin. SSR* **108** (1982) 249–252.
- [35] E. Witten, *Chern-Simons gauge theory as a string theory*, *Prog. Math.* **133** (1995) 637–678 [[hep-th/9207094](#)].
- [36] M. Herbst, C.-I. Lazaroiu and W. Lerche, *Superpotentials, $A(\infty)$ relations and WDVV equations for open topological strings*, *JHEP* **02** (2005) 071 [[hep-th/0402110](#)].
- [37] P.S. Aspinwall, *D-branes, π -stability and θ -stability*, [hep-th/0407123](#).
- [38] D.E. Diaconescu and J. Gomis, *Fractional branes and boundary states in orbifold theories*, *JHEP* **10** (2000) 001 [[hep-th/9906242](#)].